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Posted Date: 5 September 2024

doi: 10.20944/preprints202409.0460.v1

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Article

# A Duality Principle and Related Convex Dual Approximate Formulation Applied to a Non-Linear Plate Model

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**Abstract:** This article develops a duality principle applicable to originally non-convex primal variational formulations. More specifically, as a first application, we establish a convex dual approximate variational formulation for a non-linear Kirchhoff-Love plate model. The results are obtained through basic tools of functional analysis, calculus of variations, duality and optimization theory in infinite dimensional spaces. We emphasize such a final convex dual approximate formulation obtained may be applied to a large class of similar models in the calculus of variations.

**Keywords:** duality principle; non-linear plate model; convex dual approximate formulation

**MSC:** 49N15

## 1. Introduction

This article develops a duality principle applicable to a large class of models in the calculus of variations. Specifically in this text, we present applications to the non-linear Kirchhoff-Love plate model.

We emphasize the results on duality theory here addressed and developed are inspired mainly in the approaches of J.J. Telega, W.R. Bielski and co-workers presented in the articles [1–4]. Other main reference is the article by Toland, [5].

Moreover, details on the Sobolev spaces involved may be found in [6].

Similar results and models are addressed in [7–11].

Basic results on convex analysis are addressed in [12]. Other similar results and approaches may be found in [13–15].

Now we start to describe the primal variational formulation for the plate model in question.

Let  $\Omega \subset \mathbb{R}^2$  be an open, bounded and connected set with a regular (Lipschitzian) boundary denoted by  $\partial\Omega$ .

We assume such a  $\Omega$  set represents the middle surface of a thin plate with a constant thickness  $h > 0$ .

Moreover, we suppose such a plate is subject to a external load  $(P_\alpha, P) \in L^2(\Omega; \mathbb{R}^3)$  resulting a field of displacements denoted by

$$(u_\alpha, w) = (u_1, u_2, w) \in W_0^{1,2}(\Omega; \mathbb{R}^2) \times W_0^{2,2}(\Omega) = V.$$

Both the load and displacements fields refers to a cartesian system  $(0, x_1, x_2, x_3)$  and related canonical basis in  $\mathbb{R}^3$ .

Finally, we denote  $Y_1 = Y_1^* = L^2(\Omega; \mathbb{R}^4)$  and  $Y_2 = Y_2^* = L^2(\Omega; \mathbb{R}^2)$ .

We also emphasize the boundary conditions in question refer to a clamped plate.

The strain tensors are defined by

$$\gamma_{\alpha\beta}(u) = \frac{u_{\alpha,\beta} + u_{\beta,\alpha}}{2} + \frac{1}{2}w_{,\alpha}w_{,\beta},$$

and

$$\kappa_{\alpha\beta}(w) = -w_{,\alpha\beta}.$$

The plate total energy functional is defined by

$$\begin{aligned} J(u) &= \frac{1}{2} \int_{\Omega} H_{\alpha\beta\lambda\mu} \gamma_{\alpha\beta}(u) \gamma_{\lambda\mu}(u) \, dx \\ &+ \frac{1}{2} \int_{\Omega} h_{\alpha\beta\lambda\mu} \kappa_{\alpha\beta}(u) \kappa_{\lambda\mu}(u) \, dx \\ &- \langle w, P \rangle_{L^2} - \langle u_{\alpha}, P_{\alpha} \rangle_{L^2}. \end{aligned} \quad (1)$$

Here  $\{H_{\alpha\beta\lambda\mu}\}$  is a fourth order positive definite symmetric constant tensor.

Moreover

$$\{h_{\alpha\beta\lambda\mu}\} = \frac{h^2}{12} \{H_{\alpha\beta\lambda\mu}\}$$

and we denote

$$\{\bar{H}_{\alpha\beta\lambda\mu}\} = \{H_{\alpha\beta\lambda\mu}\}^{-1}$$

and

$$\{\bar{h}_{\alpha\beta\lambda\mu}\} = \{h_{\alpha\beta\lambda\mu}\}^{-1}$$

in an appropriate tensor sense.

## 2. The Main Duality Principle and Related Convex Dual Approximate Formulation

We start by defining the approximate functional  $J_1 : V \rightarrow \mathbb{R}$  by

$$J_1(u) = J(u) + \sum_{\alpha=1}^2 \frac{\varepsilon_1}{2} \int_{\Omega} (u_{\alpha})^2 \, dx,$$

and considering an appropriate real constant  $K > 0$ , the functionals  $F_1 : V \rightarrow \mathbb{R}$ ,  $F_2 : V \times Y_1^* \rightarrow \mathbb{R}$ ,  $F_3 : V \rightarrow \mathbb{R}$  and  $F_4 : V \rightarrow \mathbb{R}$ , by

$$\begin{aligned} F_1(u) &= \frac{1}{2} \int_{\Omega} h_{\alpha\beta\lambda\mu} \kappa_{\alpha\beta}(u) \kappa_{\lambda\mu}(u) \, dx + \sum_{\alpha=1}^2 \frac{K}{2} \int_{\Omega} w_{,\alpha}^2 \, dx \\ &+ \sum_{\alpha=1}^2 \varepsilon \int_{\Omega} w_{,\alpha}^2 \, dx - \langle w, P \rangle_{L^2} \end{aligned} \quad (2)$$

$$\begin{aligned} F_2(u, N) &= \frac{1}{2} \int_{\Omega} N_{\alpha\beta} w_{,\alpha} w_{,\beta} \, dx \\ &- \sum_{\alpha=1}^2 \frac{K}{2} \int_{\Omega} w_{,\alpha}^2 \, dx \end{aligned} \quad (3)$$

$$F_3(u) = \sum_{\alpha=1}^2 \frac{\varepsilon_1}{2} \int_{\Omega} u_{\alpha}^2 \, dx \quad (4)$$

$$F_4(u) = \sum_{\alpha=1}^2 \varepsilon \int_{\Omega} w_{,\alpha}^2 \, dx. \quad (5)$$

Moreover, we define the polar functionals  $F_1^* : [Y_2^*]^2 \rightarrow \mathbb{R}$ ,  $F_2^* : [Y_2^*]^2 \times Y_1^* \rightarrow \mathbb{R}$ ,  $F_3^* : Y_1^* \rightarrow \mathbb{R}$  and  $F_4^* : [Y_2^*]^2 \rightarrow \mathbb{R}$ , by

$$F_1^*(R, L) = \sup_{u \in V} \{ \langle w_{,\alpha}, R_\alpha - L_\alpha \rangle_{L^2} - F_1(u) \}, \quad (6)$$

$$\begin{aligned} F_2^*(Q, L) &= \inf_{v \in Y_2} \left\{ \langle v_\alpha, Q_\alpha + L_\alpha \rangle_{L^2} - \frac{1}{2} \int_{\Omega} N_{\alpha\beta} v_\alpha v_\beta dx \right. \\ &\quad \left. + \sum_{\alpha=1}^2 \frac{K}{2} \int_{\Omega} v_\alpha^2 dx + \frac{1}{2} \int_{\Omega} \bar{H}_{\alpha\beta\lambda\mu} N_{\alpha\beta} N_{\lambda\mu} dx \right\} \\ &= \frac{1}{2} \int_{\Omega} \overline{N_{\alpha\beta}^{(-K)}} (Q_\alpha + L_\alpha) (Q_\beta + L_\beta) dx \\ &\quad + \frac{1}{2} \int_{\Omega} \bar{H}_{\alpha\beta\lambda\mu} N_{\alpha\beta} N_{\lambda\mu} dx, \end{aligned} \quad (7)$$

if  $N = \{N_{\alpha\beta}\} \in B^*$ , where

$$B^* = \{N \in Y_1^* : \|N_{\alpha\beta}\|_\infty \leq K/8, \forall \alpha, \beta \in \{1, 2\}\},$$

$$\{N_{\alpha\beta}^{(-K)}\} = \{N_{\alpha\beta} - K\delta_{\alpha\beta}\},$$

and

$$\{\overline{N_{\alpha\beta}^{(-K)}}\} = \{N_{\alpha\beta}^{(-K)}\}^{-1}.$$

Furthermore,

$$\begin{aligned} F_3^*(N) &= \sup_{u \in V} \{ \langle u_\alpha, N_{\alpha\beta,\beta} + P_\alpha \rangle_{L^2} - F_3(u) \} \\ &= \sum_{\alpha=1}^2 \frac{1}{2\varepsilon_1} \int_{\Omega} (N_{\alpha\beta,\beta} + P_\alpha)^2 dx, \end{aligned} \quad (8)$$

$$\begin{aligned} F_4^*(R, Q) &= \sup_{(v_1, v_2) \in Y_2^*} \left\{ \langle (v_1)_\alpha, R_\alpha \rangle_{L^2} - \frac{\varepsilon}{2} \int_{\Omega} (v_1)_\alpha^2 dx \right. \\ &\quad \left. + \langle (v_2)_\alpha, Q_\alpha \rangle_{L^2} - \frac{\varepsilon}{2} \int_{\Omega} (v_2)_\alpha^2 dx \right\} \\ &= \sum_{\alpha=1}^2 \left( \frac{1}{2\varepsilon} \int_{\Omega} R_\alpha^2 dx + \frac{1}{2\varepsilon} \int_{\Omega} Q_\alpha^2 dx \right). \end{aligned} \quad (9)$$

At this point, denoting

$$D^* = \{Q = \{Q_\alpha\} \in Y_2^* : \|Q_\alpha\| \leq 5, \forall \alpha \in \{1, 2\}\},$$

we define  $J_1^* : (D^*)^2 \times B^* \times Y_2^* \rightarrow \mathbb{R}$  by

$$J_1^*(R, Q, N, L) = -F_1^*(R, L) - F_2^*(Q, L, N) - F_3^*(N) + F_4^*(R, Q),$$

and  $J_2^* : (D^*)^2 \times B^* \rightarrow \mathbb{R}$  by

$$J_2^*(R, Q, N) = \text{sta}_{L \in Y_2^*} J_1^*(R, Q, N, L) = J_1^*(R, Q, N, \hat{L}(R, Q, N)),$$

where  $\hat{L} = \hat{L}(R, Q, N) \in Y_2^*$  is the only solution of the linear equation in  $L$

$$\frac{\partial J_1^*(R, Q, N, \hat{L})}{\partial L} = \mathbf{0}.$$

Moreover, we define  $J_3^* : (D^*)^2 \times B^* \rightarrow \mathbb{R}$ , by

$$\begin{aligned} & J_3^*(R, Q, N) \\ = & J_2^*(R, Q, N) \\ & - \sum_{\alpha=1}^2 \sum_{\beta=1}^2 \frac{K_1}{2} \left\| \bar{H}_{\alpha\beta\lambda\mu} N_{\lambda\mu} - \frac{(N_{\alpha\rho\rho} + P_\alpha)_{,\beta} + (N_{\beta\rho\rho} + P_\alpha)_{,\alpha}}{2\varepsilon_1} - \frac{1}{2} \tilde{v}_\alpha \tilde{v}_\beta \right\|_{0,2}^2, \end{aligned} \quad (10)$$

where

$$\tilde{v}_\alpha = \overline{N_{\alpha\beta}^{(-K)}}(Q_\beta + L_\beta(R, Q, N)), \quad \forall \alpha \in \{1, 2\}.$$

Here, we assume

$$K_1 \gg \max\{1, K, \max\{\bar{h}_{\alpha\beta\lambda\mu}, \alpha, \beta, \lambda, \mu \in \{1, 2\}\}\},$$

$$0 < \varepsilon, \varepsilon_1 \ll 1$$

and

$$\frac{1}{\varepsilon} \gg \max\left\{K_1, \frac{1}{\varepsilon_1}\right\}.$$

Observe that

$$\frac{\partial^2 J_3^*(R, Q, N)}{\partial Q_\alpha^2} = \mathcal{O}\left(\frac{1}{\varepsilon}\right) > \mathbf{0},$$

$$\frac{\partial^2 J_3^*(R, Q, N)}{\partial R_\alpha^2} = \mathcal{O}\left(\frac{1}{\varepsilon}\right) > \mathbf{0},$$

$\forall \alpha \in \{1, 2\}$ .

Thus, considering also the remaining mixed variations in  $Q_\alpha$  and  $R_\alpha$ , we may infer that

$$\det \left\{ \frac{\partial^2 J_3^*(R, Q, N)}{\partial Q_\alpha \partial R_\beta} \right\} > \mathbf{0},$$

in  $(D^*)^2 \times B^*$ .

Moreover, by direct computation, clearly

$$\det \left\{ \frac{\partial^2 J_3^*(R, Q, N)}{\partial N_{\alpha\beta} \partial N_{\lambda\mu}} \right\} < \mathbf{0},$$

in  $(D^*)^2 \times B^*$ .

From such results, we may infer that  $J_3^*$  is convex in  $(R, Q)$  and concave in  $N$  in  $(D^*)^2 \times B^*$ .

Let  $(\hat{R}, \hat{Q}, \hat{N}, \hat{L}) \in (D^*)^2 \times B^* \times Y_2^*$  be such that

$$\delta J_1^*(\hat{R}, \hat{Q}, \hat{N}, \hat{L}) = \mathbf{0}.$$

Let  $u_0 = ((u_0)_\alpha, w_0) \in V$  be such that

$$(u_0)_\alpha = \frac{\hat{N}_{\alpha\beta,\beta} + P_\alpha}{\varepsilon_1},$$

and

$$(w_0)_{,\alpha} = \frac{\hat{Q}_\alpha}{\varepsilon},$$

$\forall \alpha \in \{1, 2\}$ .

From standard results in Duality Theory and the Legendre Transform properties, we may obtain

$$\delta J_1(u_0) = \mathbf{0},$$

$$\delta J_2^*(\hat{R}, \hat{Q}, \hat{N}) = \mathbf{0}$$

,

$$\delta J_3^*(\hat{R}, \hat{Q}, \hat{N}) = \mathbf{0}$$

and

$$\begin{aligned} J_1(u_0) &= J_1^*(\hat{R}, \hat{Q}, \hat{N}, \hat{L}) \\ &= J_2^*(\hat{R}, \hat{Q}, \hat{N}) \\ &= J_3^*(\hat{R}, \hat{Q}, \hat{N}). \end{aligned} \tag{11}$$

From such results and the Min-Max Theorem, we have

$$J_3^*(\hat{R}, \hat{Q}, \hat{N}) = \inf_{(R, Q) \in (D^*)^2} \left\{ \sup_{N \in B^*} J_3^*(R, Q, N) \right\}.$$

Joining the pieces, we have got

$$\begin{aligned} J_1(u_0) &= J_1^*(\hat{R}, \hat{Q}, \hat{N}, \hat{L}) \\ &= J_2^*(\hat{R}, \hat{Q}, \hat{N}) \\ &= \inf_{(R, Q) \in (D^*)^2} \left\{ \sup_{N \in B^*} J_3^*(R, Q, N) \right\} \\ &= J_3^*(\hat{R}, \hat{Q}, \hat{N}). \end{aligned} \tag{12}$$

**Remark 2.1.** Defining  $J_5^* : (D^*)^2 \rightarrow \mathbb{R}$  by

$$J_5^*(R, Q) = \sup_{N \in B^*} J_3^*(R, Q, N),$$

we have that such a functional  $J_5^*$  is convex in  $(D^*)^2$  as the supremum in  $N \in B^*$  of a family of convex functionals in  $(R, Q)$ .

In such a case, we have also obtained

$$\begin{aligned} J_1(u_0) &= J_1^*(\hat{R}, \hat{Q}, \hat{N}, \hat{L}) \\ &= J_2^*(\hat{R}, \hat{Q}, \hat{N}) \\ &= \inf_{(R, Q) \in (D^*)^2} \left\{ \sup_{N \in B^*} J_3^*(R, Q, N) \right\} \\ &= J_3^*(\hat{R}, \hat{Q}, \hat{N}) \\ &= \inf_{(R, Q) \in (D^*)^2} J_5^*(R, Q) \\ &= J_5^*(\hat{R}, \hat{Q}). \end{aligned} \tag{13}$$

### 3. Conclusion

In this article, we have developed a duality principle and related convex dual approximate variational formulation for an originally non-convex primal one.

We highlight the results here obtained are applicable to a large class of models in the calculus of variations, including other plate and shell non-linear theories, models in superconductivity, phase transition and micro-magnetism, among many others.

In a near future research we intend to apply such results to some of these mentioned related models.

**Conflicts of Interest:** The author declares no conflict of interest concerning this article.

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